



TITLE:

# Antiferromagnetism of $S=1/2$ Triangles in $\text{La}_4\text{Cu}_3\text{MoO}_{12}$ (SOLID STATE CHEMISTRY- Multicomponent Materials)

AUTHOR(S):

Azuma, Masaki; Ishiwata, Shintaro; Takano, Mikio

---

CITATION:

Azuma, Masaki ...[et al]. Antiferromagnetism of  $S=1/2$  Triangles in  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$  (SOLID STATE CHEMISTRY-Multicomponent Materials). ICR Annual Report 2001, 7: 20-21

ISSUE DATE:

2001-03

URL:

<http://hdl.handle.net/2433/65280>

RIGHT:

# Antiferromagnetism of $S=1/2$ triangles in $\text{La}_4\text{Cu}_3\text{MoO}_{12}$

Masaki Azuma, Shintaro Ishiwata and Mikio Takano

Magnetic properties of a cupric triangular cluster compound  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$  were investigated. Susceptibility data show that paramagnetic cupric spin ( $S=1/2$ ) above room temperature form  $S=1/2$  trimers at lower temperatures because of strong intra-trimer antiferromagnetic interactions. The cluster moments acquire antiferromagnetic order below 2.6 K. The trimers can be polarized in an external field of 20 T at 1.3 K. The magnetization remains nearly constant at  $1 \mu_B$  per trimer up to a magnetization plateau at 55 T.

**Keywords:** Spin trimer/ Magnetic susceptibility/ High field magnetization/ Neutron scattering

The spin quantum number (small, large, integer, half-integer) and the geometry of the magnetic sublattice affect the ground state of quantum antiferromagnets in various fashions. For examples, spin singlet ground states with finite gaps to magnetic excited states have been found in several 1D antiferromagnetic (AF) systems such as  $S=1/2$  alternating bond chains,  $S=1/2$  2-leg ladders and  $S=1$  chains (Haldane systems). The discovery of inorganic model compounds such as  $\text{CuGeO}_3$  (spin-Peierls),  $(\text{VO})_2\text{P}_2\text{O}_7$  (alternating chain) [1],  $\text{Y}_2\text{BaNiO}_5$  (Haldane),  $\text{SrCu}_2\text{O}_3$  [2] and  $\text{Sr}_{14}\text{Cu}_{24}\text{O}_{41}$  (ladders) in the past decade has stimulated keen interest in such 1D systems. Their large AF interactions ( $J$ ) and the thermal stability make it rather easy to finely tune the electronic properties of these materials by means of chemical doping.

In the gapped ground states of these compounds, two neighboring  $S=1/2$  spins form a spin singlet. On the other hand, AF trimer compounds are rare in nature.  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$  is a layered compound with an average structure of the  $\text{YAlO}_3$  type [3]. The  $\text{Cu}_3\text{MoO}_4$  layer of this compound shown in Fig. 1 can be derived from a triangular  $\text{CuO}$  layer by replacing a quarter of the  $\text{Cu}^{2+}$  ions with nonmagnetic  $\text{Mo}^{6+}$  ions. One can consider the

$\text{Cu}_3\text{MoO}_4$  layer as being made of  $\text{Cu}_3\text{O}$  triangular clusters as suggested from the bond lengths: The average Cu-O bond lengths within the triangle,  $1.983 \text{ \AA}$ , is much shorter than that between neighboring triangles of  $2.690 \text{ \AA}$ . The  $\text{Cu}_3\text{MoO}_4$  layers are separated from each other by  $/\text{O}/\text{La}/\text{O}/$  layers. The structure can be described as a quasi-2D orthorhombic lattice made of slightly distorted cupric ( $S=1/2$ ) trimers [4].

Figure 2 shows the temperature dependences of magnetic susceptibility and its inverse measured on heating from 5 to 800 K in an external field of 1 T. The susceptibility above 400 K was fitted well to the Curie-Weiss law with  $\mu_{\text{eff}} = 1.81$  ( $g = 2.09$ ) and a Weiss temperature ( $\theta$ ) of  $-558 \text{ K}$ . On the other hand, the slope of the  $1/\chi$ - $T$  plot below 250 K decreased to 0.39 times the high-temperature value. The small Curie constant ( $C$ ) below 250 K indicates that each trimer has a total spin  $S_{\text{total}}=1/2$ . The large intra-trimer AF interactions give the large Weiss constant at high temperatures. The inter-trimer interactions seem to be weakly AF as suggested from the small Weiss constant of  $-16 \text{ K}$  derived from the low temperature data.

In the data measured in a field of 0.1 T down to 1.8 K shown in the inset of Fig. 2, a maximum appeared at 5 K

## SOLID STATE CHEMISTRY — Multicomponent Materials —

### Scope of research

*Novel 3d transition-metal oxides showing exotic electrical and magnetic properties are being searched for using different synthesizing techniques like high pressure synthesis (5 GPa and 1000°C, typically) and epitaxial film growth. Recent topics are:*

- High  $T_c$  superconductivity.
- Low-dimensional spin systems like ladders showing dramatic quantum effects.
- Oxides of late 3d transition metals like  $\text{SrFeO}_3$  with strong oxygen-hole character.



Prof  
TAKANO, Mikio  
(D Sc)



Assoc Prof  
TERASHIMA, Takahito  
(D Sc)



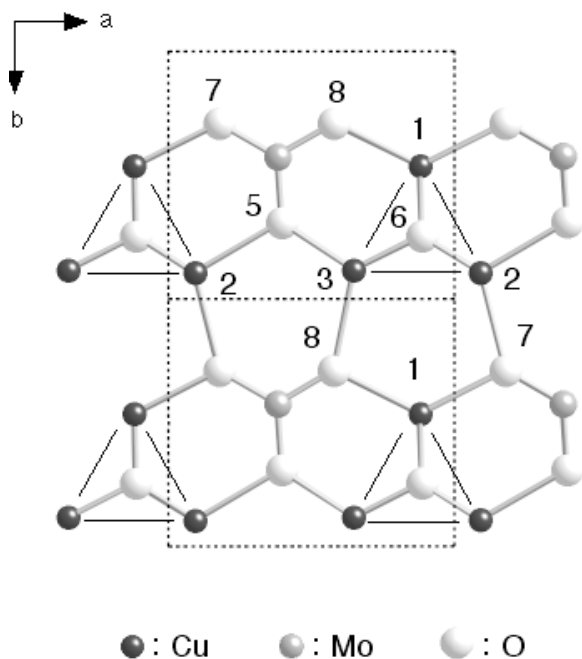
Instr  
AZUMA, Masaki  
(D Sc)

### Students:

FURUBAYASHI, Yutaka (DC)  
HAYASHI, Naoaki (DC)  
SAITO, Takashi (DC)  
ISHIWATA, Shintaro (MC)  
MASUNO, Atsunobu (MC)  
YAMAMOTO, Masamichi (MC)  
YOSHIDA, Hirohumi (MC)  
RIJSSENBEEK, Job T. (RS)  
Ninjabdar Tsedev (RS)

### Research Fellow:

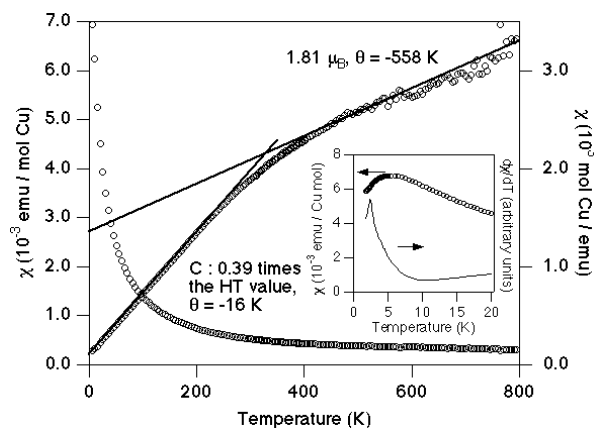
ICHIKAWA, Noriya  
KAWASAKI, Shuji



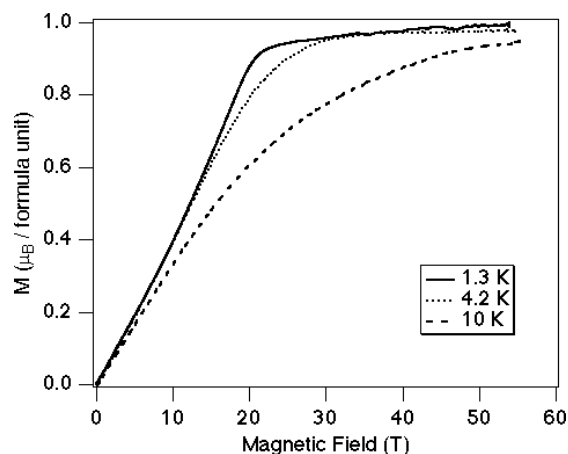
**Figure 1.**  $\text{Cu}_3\text{MoO}_4$  plane of  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$ . The numbers show the crystallographic sites. The solid line represent the triangle clusters.

and a kink at 2.6 K followed. The existence of the kink was clearly seen as a peak in the specific heat data. As can be seen in Fig. 1 there is no spin frustration among the  $S_{\text{total}}=1/2$  spins, so it is reasonable to assume that the anomaly results from AF order of spin trimers. The susceptibility maximum at 5 K can be attributed to short range order in the two dimensional spin system. The elastic neutron scattering study confirmed the antiferromagnetic order with an ordered moment of  $0.8 \mu_B$  per trimer.

We propose the following picture for the present compound. Above 400 K, the susceptibility could be explained assuming  $S=1/2$  spin localized on each Cu ion. The large AF interactions within the triangles led to a large Weiss temperature of -558 K. With decreasing the



**Figure 2.** Temperature dependence of magnetic susceptibility of  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$  below 800 K. The solid lines correspond the fit to the Curie-Weiss law. The inset shows the data taken at 0.1 T.



**Figure 3.** Magnetization curve of  $\text{La}_4\text{Cu}_3\text{MoO}_{12}$  measured in pulsed magnetic field at 1.3, 4.2 and 10 K.

temperature, the Curie constant decreased to about 1/3 of the high-temperature value because only the ground state with  $S_{\text{total}}=1/2$  per cluster was populated. Since the inter-triangle interactions are weak, the Weiss temperature is as small as -16 K. The  $S_{\text{total}}=1/2$  spins localized on the trimers order antiferromagnetically at 2.6 K. If inter-triangle interactions are weak, it should be possible to align the ordered spins by applying magnetic field (spin flip). The results of the magnetization measurements at 1.3, 4.2 and 10 K are shown in Fig. 3. At 1.3 K, the data below 20 T were concave reflecting the AF ordering, whereas above 20 T the magnetization tended to saturate at  $1 \mu_B$  / formula unit, *i. e.*,  $1 \mu_B$  per trimer. This confirms that the system orders antiferromagnetically keeping the nature of the ground state with  $S_{\text{total}}=1/2$ . The full magnetization of this system should be  $3 \mu_B$  / formula unit, so the observed saturation-like behavior is actually an intermediate plateau. The excitation energy from  $S_{\text{total}}=1/2$  to  $S_{\text{total}}=3/2$  state is  $3/2 J$  for an equilateral triangle. Applying  $J=813$  K estimated from the fitting of the susceptibility data, it is expected that another spin flip will be observed at 1500 T. Actually, inelastic neutron scattering data showed a well defined excitation at  $\Delta E=133$  meV (1542 K) which can be associated with an intratrimer doublet-quartet ( $S_{\text{total}}=1/2$  to  $S_{\text{total}}=3/2$ ) transition.

## Reference

1. M. Azuma, T. Saito, Y. Fujishiro, Z. Hiroi, M. Takano, F. Izumi, T. Kamiyama, T. Ikeda, Y. Narumi and K. Kindo, Phys. Rev. B **60**, 10145 (1999).
2. Z. Hiroi, M. Azuma, M. Takano and Y. Bando, J. Solid State Chem. **95**, 230 (1991); M. Azuma, Z. Hiroi, M. Takano, K. Ishida and Y. Kitaoka, Phys. Rev. Lett. **73**, 3463 (1994).
3. D. A. Vander Griend, S. Boudin, V. Caignaert, K. R. Poeppelmerer, Y. Wang, V. P. Dravid, M. Azuma, M. Takano, Z. Hu and J. D. Jorgensen, J. Am. Chem. Soc. **121**, 4787 (1999).
4. M. Azuma, T. Odaka, M. Takano, D. A. VanderGriend, K. R. Poeppelmeier, Y. Narumi, K. Kindo, Y. Mizuno and S. Maekawa, Phys. Rev. B **62**, R3588 (2000).